

# Forward-backward correlations and charged-particle azimuthal distributions in $pp$ interactions using the ATLAS detector

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## The ATLAS collaboration

**ABSTRACT:** Using inelastic proton-proton interactions at  $\sqrt{s} = 900 \text{ GeV}$  and  $7 \text{ TeV}$ , recorded by the ATLAS detector at the LHC, measurements have been made of the correlations between forward and backward charged-particle multiplicities and, for the first time, between forward and backward charged-particle summed transverse momentum. In addition, jet-like structure in the events is studied by means of azimuthal distributions of charged particles relative to the charged particle with highest transverse momentum in a selected kinematic region of the event. The results are compared with predictions from tunes of the PYTHIA and HERWIG++ Monte Carlo generators, which in most cases are found to provide a reasonable description of the data.

**KEYWORDS:** Hadron-Hadron Scattering

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**1 Introduction**

Two general classes of quantum chromodynamic processes combine to generate the events that are seen in high energy hadron collisions. The first of these comprises soft processes, which are commonly pictured in terms of multiple parton interactions at low transverse momentum,  $p_T$ , modelled on multiple  $t$ -channel gluon exchanges. Processes of this kind, together with non-perturbatively described contributions such as the formation of beam remnants, are characterised by correlations between numbers of particles per unit of pseudorapidity over an extended pseudorapidity range [1]. At higher  $p_T$  values, harder processes arise that are dominated by perturbatively-described radiation or partonic scatters and are produced by single- and few-parton exchanges. These can generate particle jets, which are by definition characterised by short-range correlation; here high particle multiplicity is

found within one unit in pseudorapidity, quickly diminishing as the pseudorapidity distance increases. Thus, soft processes produce lower particle multiplicities with weaker correlations over a wider pseudorapidity distance, so-called “long-range correlations”, in comparison to the processes where a small number of jets are formed. As particle  $p_T$  values increase above a few hundred MeV, there is a gradual transition between the soft regime and the hard. Hence, the correlation measured as a function of separation in pseudorapidity and of  $p_T$  probes the strength of each contribution.

A number of different event generators are available that combine the basic physical processes appropriately [2]; however, the more detailed aspects of the models need to be determined from data, a process referred to as “tuning” the models. For example, there are choices to be made as to whether the initial- and final-state gluon radiation are interleaved with the multiple parton interactions (MPI) and whether the modelling of the showers is to be ordered in terms of  $p_T$ , angle, or the momentum transfer squared,  $Q^2$ , of the parton scatter. String or cluster fragmentation may be selected, and diffractive processes should be taken into consideration. Factors such as these can affect the correlations between the generated particles.

The tunings of the non-diffractive models have made use of output from many experiments; a variety of event characteristics were taken into account, such as distributions in particle multiplicity, transverse momentum and rapidity. A number of different tunes of the leading-order plus leading-logarithm generator PYTHIA [3–5] exist that provide a reasonably good description of events up to LHC energies [6–8]. The PYTHIA model uses string fragmentation [9, 10], in contrast to the cluster fragmentation used in the HERWIG model [11–13].

The correlations in pseudorapidity between the particles in an event have not been explicitly considered in the existing Monte Carlo (MC) tunes, and provide a further means of improving and helping to discriminate between the models [14]. Such correlations are the subject of the first part of the present study, which was carried out using proton-proton collisions recorded with the ATLAS detector [15] at the LHC.<sup>1</sup> The measurements were inspired by earlier work by the UA5 experiment [16] and follow published studies at the LHC of the multiplicities themselves [17–21]. They also complement an earlier analysis by CMS of charged-particle correlations [22] and a recent analysis by ATLAS of two-particle correlations [23]. The approach taken here is to choose pairs of pseudorapidity intervals equal in size and symmetrically located in the forward and backward direction, allowing a measurement of what will be termed a forward-backward (FB) correlation between quantities in the two intervals. To determine the effects of both short-range and long-range physical correlations, the FB correlations are measured for increasing separation between the two selected pseudorapidity intervals.

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<sup>1</sup>ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the  $z$ -axis along the beam pipe. The  $x$ -axis points from the IP to the centre of the LHC ring, and the  $y$ -axis points upward. Cylindrical coordinates  $(r, \phi)$  are used in the transverse plane,  $\phi$  being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle  $\theta$ , with respect to the  $z$ -axis, as  $\eta = -\ln \tan(\theta/2)$ . Throughout the paper the term “forward” is used to refer to the direction of positive  $z$  or  $\eta$ , and “backward” to refer to the negative direction.

Two types of FB correlation are presented: correlations in charged-particle multiplicity and correlations in the summed values of the transverse momenta of the charged particles. A further perspective comes from increasing the value of the minimum transverse momentum,  $p_{T\min}$ , of the particles that are considered. In this way the effects of the soft processes can be gradually reduced, while the effects of the harder processes remain. The use of data taken at two very different proton-proton centre-of-mass energies,  $\sqrt{s} = 900$  GeV and 7 TeV, gives further information that can be compared with MC tunes based on  $p\bar{p}$  interactions at the intermediate energy of the Tevatron, at  $\sqrt{s} = 1.96$  TeV. The 7 TeV data provide access to a new energy regime, which is the main goal of the LHC. It is important to understand the relationship between softer and harder processes in these two regimes in order to confirm our understanding of the background processes at the energies where new physics might be found.

Although such correlation measurements are sensitive to the presence of jets in events, no jet features are explicitly sought or exploited. Jets are by definition associated with short-range correlations in both pseudorapidity and azimuthal angle, and to obtain additional information about these correlations, a further analysis was performed using a different technique. Since the direction of a jet is often close to that of the highest- $p_T$  particle in the event, the charged particle with the highest  $p_T$  is identified and the distributions of the other charged particles are studied as a function of their azimuthal angle relative to its direction.

The distribution of particles in the azimuthal plane of an event has been studied before at Tevatron energies [24–26] and at the LHC [27]. Taken relative to the highest- $p_T$  particle, the azimuthal distribution of particles shows two peaks, centred at zero and at  $\pi$ , that can indicate jet structure. In the events studied in the present measurement, however, these peaks are typically not observed in single events, most of which contain too few hard particles to form clearly defined jets, but become visible statistically when assembling the distribution over many events. The particle distributions also display the so-called “underlying event”, the effects of which are particularly apparent at azimuthal angles perpendicular to the jets [27]. However, the principal goal of the present study is an understanding of the shape of the first jet-like peak, which is described differently depending on the MC model used. Again, it is found that the transition from 900 GeV to 7 TeV LHC energies provides a significantly different perspective on the physics.

This paper is structured as follows. The statistical methods for the correlation analysis and the azimuthal analysis are described in sections 2 and 3. Section 4 describes the ATLAS detector and the event selection, while the MC samples are described in section 5. Details of the analysis method are given in section 6. The systematic uncertainties are discussed in section 7, and the results and conclusions are presented in sections 8 and 9.

## 2 Forward-backward correlations

All the correlations measured in this study, whether from symmetric or asymmetric pseudorapidity intervals, are referred to as forward-backward (FB) correlations and are measured using charged particles only. The FB multiplicity correlation,  $\rho_{fb}^n$ , between two particle

multiplicities is the normalised covariance between the two distributions, relative to the mean value of each. For a given data sample or MC sample, it is defined in the usual way as:

$$\rho_{fb}^n = \frac{\langle (n_f - \langle n_f \rangle)(n_b - \langle n_b \rangle) \rangle}{\sqrt{\langle (n_f - \langle n_f \rangle)^2 \rangle \langle (n_b - \langle n_b \rangle)^2 \rangle}} = \frac{\sum x_f^n x_b^n}{N \sigma_f^n \sigma_b^n}. \quad (2.1)$$

Here,  $n_f$  and  $n_b$  are the respective multiplicities of particles of interest in two chosen forward and backward intervals in an event,  $\langle \rangle$  denotes a mean over the events in the sample,  $N$  is the total number of events,  $x_f^n, x_b^n$  are the differences between  $n_f, n_b$  and their means, and  $\sigma_f^n, \sigma_b^n$  are the standard deviations of  $n_f$  and  $n_b$  about their means. The sum in the right-hand side of the equation is taken over the events in the sample. In the present measurement, particles of interest are those above a given  $p_{T\min}$  value. Intervals in pseudorapidity of size  $\delta\eta = 0.5$  were chosen, which allows good statistical accuracy on each measured point while preserving a sensitivity to physically interesting variations with pseudorapidity  $\eta$ . A range  $-2.5 < \eta < 2.5$  was considered, corresponding to the inner tracking detector acceptance, divided into five forward and five backward intervals.

The FB momentum correlation,  $\rho_{fb}^{p_T}$ , between two summed transverse momentum ( $\sum p_T$ ) values is defined similarly:

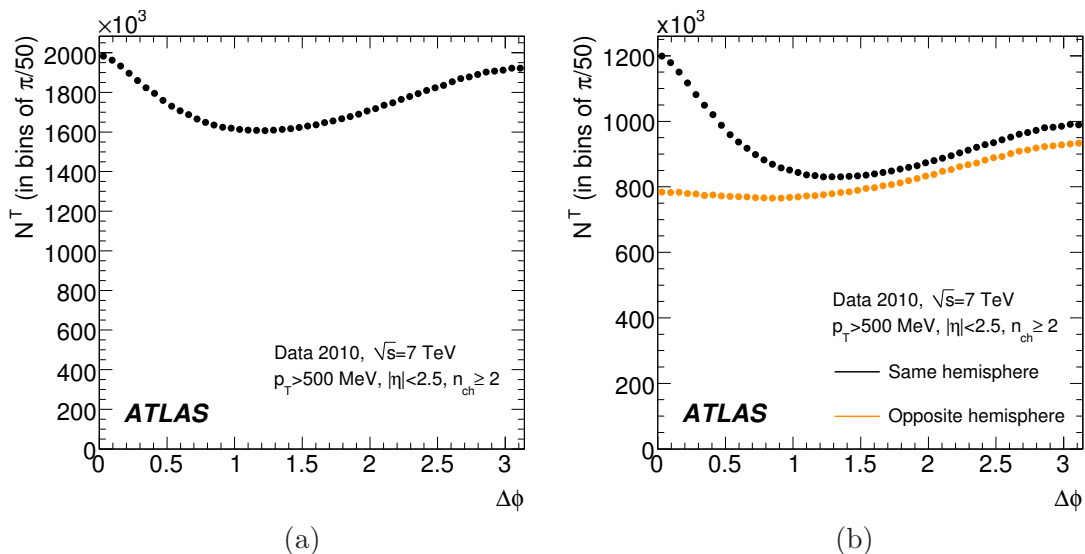
$$\rho_{fb}^{p_T} = \frac{\langle (\sum p_T^f - \langle \sum p_T^f \rangle)(\sum p_T^b - \langle \sum p_T^b \rangle) \rangle}{\sqrt{\langle (\sum p_T^f - \langle \sum p_T^f \rangle)^2 \rangle \langle (\sum p_T^b - \langle \sum p_T^b \rangle)^2 \rangle}} = \frac{\sum x_f^{p_T} x_b^{p_T}}{N \sigma_f^{p_T} \sigma_b^{p_T}}, \quad (2.2)$$

where  $\sum p_T^f$  and  $\sum p_T^b$  are the sums of the absolute transverse momentum values of the charged particles in the two chosen forward and backward intervals in an event,  $x_f^{p_T}, x_b^{p_T}$  are the differences between  $\sum p_T^f, \sum p_T^b$  and their means, and  $\sigma_f^{p_T}, \sigma_b^{p_T}$  are the respective standard deviations. The  $p_{T\min}$  values used in measuring the correlations are discussed in the sections below.

### 3 Azimuthal distributions

To obtain sensitivity to the presence of jets and jet-like structures in the events, the highest- $p_T$  charged particle in a given  $|\eta|$  region in an event is identified and termed the ‘‘leading particle’’. The azimuthal difference  $\Delta\phi$  between the leading particle and any other accepted charged particle is the unsigned angle (in the range  $0-\pi$ ) between the two particles in the transverse plane. The variable  $N^T$  is defined as the number of selected charged particles, taken over the entire event sample, in a given interval of  $\Delta\phi$  of width  $\delta\phi$ , where  $\delta\phi$  takes the value  $\pi/50$  in the present analysis. A  $p_T > p_{T\min}$  requirement is again imposed on the selected particles.

In the general case of jet production, it is expected that the distribution in  $N^T$  should show peaks at  $\Delta\phi = 0$  and at  $\Delta\phi = \pi$ , with a minimum lying in between. As observed in the present study (figure 1a), this minimum takes the form of a flat contribution, or pedestal, which arises from the overall charged-particle activity in the event [27] together with uncorrelated backgrounds that are hard to model, such as fake tracks in the detector. The present study focuses on the peak structure, suppressing the influence of the pedestal



**Figure 1.** Azimuthal distributions for tracks with  $p_{T\min} = 500$  MeV and  $|\eta| < 2.5$  in data at  $\sqrt{s} = 7$  TeV. (a) Uncorrected distribution of  $\Delta\phi$ ; (b) distributions for the same hemisphere in  $\eta$  as the leading track, and for the opposite hemisphere. The statistical uncertainties are smaller than the plotted symbol.

on the measurement by performing a subtraction. Since the pedestal subtraction is performed in a corresponding way on the MC models tested, the latter do not need to model the pedestal level accurately for this analysis. In each case, we fit the observed  $N^T$  distribution with a quadratic polynomial, determine the minimum value  $N_{\min}^T$ , and subtract this value from all the  $N^T$  values in the distribution. The “subtracted” number of charged particles per radian in azimuthal angle is then normalised to give the quantity

$$N_{\text{sub}}^T = \frac{(N^T - N_{\min}^T)}{\sum(N^T - N_{\min}^T)}, \tag{3.1}$$

where the sum is taken over the bins of  $\delta\phi$  used in the analysis.

Averaged over many events, the  $N_{\text{sub}}^T$  variable is sensitive to the presence of jet pairs whenever both jets occur within the accepted  $\eta$  interval. The direct effect of the second jet is removed, suppressing pedestal effects in an alternative way, by examining a second variable. We distinguish between particles produced in the same forward-backward event hemisphere in  $\eta$  as the leading particle, and those produced in the opposite hemisphere ( $\eta * \eta_{\text{lead}} \geq 0$  and  $< 0$ , respectively), denoting these as  $N_{\text{same}}^T$  and  $N_{\text{opp}}^T$  respectively, as shown in figure 1b. The normalised asymmetry is then defined as

$$N_{\text{SO}}^T = \frac{(N_{\text{same}}^T - N_{\text{opp}}^T)}{\sum(N_{\text{same}}^T - N_{\text{opp}}^T)}. \tag{3.2}$$

On the assumption that the leading particle indicates the possible presence of a jet,  $N_{\text{SO}}^T$  is expected to take large positive values for  $\Delta\phi \approx 0$  but to be smaller for larger  $\Delta\phi$ . Thus the jet structure associated with the leading particle is well-characterised, but the effects

of the opposing jets tend on average to cancel out. Any non-zero value at large  $\Delta\phi$  reflects long-range correlations in  $\eta$  with jet effects largely removed. In this way, the two variables enhance the sensitivity to different aspects of jet-like events, and are sensitive to an aspect of long-range correlations that complements the FB measurements.

## 4 Detector and event selection

The ATLAS detector at the LHC covers almost the entire solid angle around the collision point with layers of tracking detectors, calorimeters and muon chambers. In the present analysis, tracks and vertices are reconstructed using the inner detector, which consists of a silicon pixel detector, a silicon strip detector and a transition radiation tracker. These extend in total from radii of 50.5 mm to 1066 mm from the beam line, and are located in a 2 tesla axial magnetic field provided by a superconducting solenoid. A track from a charged particle has typically 11 silicon detector hits and more than 30 transition radiation tracker hits. Coverage is obtained in the pseudorapidity range  $|\eta| < 2.5$ .

The data analysed here were collected with the ATLAS detector in 2009 at  $\sqrt{s} = 900$  GeV, and in early 2010 at  $\sqrt{s} = 7$  TeV, corresponding to integrated luminosities of  $7 \mu\text{b}^{-1}$  and  $190 \mu\text{b}^{-1}$  respectively. The maximum instantaneous luminosity was  $\sim 1.9 \times 10^{27} \text{ cm}^{-2}\text{s}^{-1}$  (at 7 TeV), and the probability of additional interactions in the same bunch crossing as a given event was of the order of  $10^{-3}$ .

This analysis is based on the same dataset as was used for the ATLAS minimum-bias particle multiplicity analysis [17] and uses the same minimum-bias scintillator trigger, the same event selections, and the same quality criteria on the track reconstruction and selection. An event in this dataset was required to have at least one primary vertex, formed from a minimum of two tracks with  $p_T > 100$  MeV and consistent with the beamline. The primary vertex with tracks giving the highest  $\Sigma p_T^2$  was selected, and events with a second primary vertex with four or more tracks were rejected, to minimise any possible effects due to pile-up. Tracks with an impact parameter above 1.5 mm in the transverse plane were removed to reduce the effects of particles from secondary decays. Around ten million events at 7 TeV and one million events at 900 GeV were selected at this stage.

For the FB correlation analysis, event samples were formed having at least two tracks in the event selected with  $p_T > p_{T\text{min}}$  and  $|\eta| < 2.5$ , for a range of  $p_{T\text{min}}$  values: 100, 200, 500 MeV, 1.0, 1.5, 2.0 GeV (see section 8.1). For the azimuthal analysis, a value of  $p_{T\text{min}} = 500$  MeV was used in the final track and event selection, as discussed in section 6.2.

## 5 Monte Carlo models

It was necessary to apply corrections to the observed FB correlations in order to obtain the final results. These corrections were evaluated using simulated events generated with the MC09 tune of PYTHIA 6 [4, 28]. This tune, based on PYTHIA 6.4.21, was created by ATLAS to emulate the particle production at the Tevatron and provide predictions for the LHC. In PYTHIA 6, the initial-state gluon radiation is interleaved with the MPI, which is expected to reduce the difference in strength between short- and long-range correlations.

Four further tunes, chosen to make use of different physical approaches as described below, were employed to evaluate model-dependent systematic effects. Around 800,000 MC events for each tune were used. Each of these samples contained a mixture of non-diffractive, single-diffractive and double-diffractive events. The subprocesses were combined according to the generator cross-sections. All the MC samples were passed through standard ATLAS software programs [29], based on GEANT4 [30], to simulate the trigger and detector effects before event reconstruction.

Two of the further tunes were also based on PYTHIA 6.4.21. The DW tune [31] was created to describe underlying-event (UE) data at the Tevatron and uses a  $Q^2$ -ordered shower, in contrast to the  $p_T$  shower ordering used by the other PYTHIA tunes. In the Perugia0 [32] tune the soft contribution was optimised using minimum-bias data from the Tevatron and the CERN Sp $\bar{p}$ S colliders. A PYTHIA 8 sample [5] was also used, with default parameter values (Tune 1). In PYTHIA 8, the final-state radiation (FSR) as well as the initial-state radiation (ISR) are interleaved with the MPI [33]. An alternative approach was provided by using PHOJET 1.12.3.15 [34]. This model uses a dual-parton event description, with soft hadronic processes described by pomeron exchange and semi-hard processes described by perturbative parton scattering. It has similar fragmentation and hadronisation to the PYTHIA model, and includes an additional central diffractive component.

The final results are compared with a selection of the above tunes and additional tunes, some of which have made use of LHC minimum-bias and underlying-event data. None of the tunes has taken account of FB correlation measurements. In general, the effect of the newer tunes has been to change the numbers of multiple parton interactions and the degree of colour reconnection. Perugia2011 [32] is a more recent tune of PYTHIA 6.4 that uses the usual  $p_T$ -ordered shower model, the interleaved MPI model, and a new set of colour reconnections. AMBT2B is a tune of PYTHIA 6.4 to ATLAS minimum-bias data [35]. The new PYTHIA 8 4C tune was used [5], and also HERWIG++ [12, 13], in versions 2.1 and 2.5, tuned to minimum-bias LHC data at 900 GeV (tune MU900-2) and to the underlying event at 7 TeV (tune UE7-2). The HERWIG model uses an angular-ordered parton shower and cluster fragmentation, but does not include diffractive processes, which constitute approximately 35% of the PYTHIA cross section and 20% of that of PHOJET. Appendix A presents in tabular form the main features of the PYTHIA and HERWIG tunes that have been used.

## 6 Analysis method

The correlations and distributions of physical interest are defined at the so-called “particle level” of the events; that is to say, in terms of the stable charged particles produced in the primary interaction with lifetime  $c\tau > 10$  mm. The kinematic region in a given measurement was defined by a requirement that there should be at least two such particles associated with the event, having  $p_T > p_{T\min}$  and  $|\eta|$  within a specified range up to a maximum value of 2.5. The FB correlations and the azimuthal distributions observed in the detector were corrected to the applicable kinematic region at the particle level.



### 6.1 Forward-backward correlations

As observed in the detector, the FB correlations are modified by detector effects such as track reconstruction efficiency, trigger efficiency, vertex finding efficiency, and the detection of secondary particles arising from decays and interactions in the detector material. A simple weighting-factor correction to the number of measured tracks in a given interval (or to their summed  $p_T$ ) will not suffice when correlations are being measured, because the corrections may themselves be correlated between the different intervals. In the context of a set of variables with multiple correlations between them, a natural method to use for determining the correlations is multiple regression theory [36]. The following multiple regression technique was used to apply corrections to the observed FB correlations so as to obtain the FB correlations at the particle level.

In general, given four statistically varying quantities that may be correlated among themselves, the linear regression relationships take the form

$$\begin{aligned}
 x_1 &= b_{12.34} x_2 + b_{13.24} x_3 + b_{14.23} x_4 \\
 x_2 &= b_{21.34} x_1 + b_{23.14} x_3 + b_{24.13} x_4 \\
 x_3 &= b_{31.24} x_1 + b_{32.14} x_2 + b_{34.12} x_4 \\
 x_4 &= b_{41.23} x_1 + b_{42.13} x_2 + b_{43.12} x_3.
 \end{aligned}
 \tag{6.1}$$

Each variable  $x_i$  denotes the difference between a statistically varying quantity and its sample mean. In the present case, this notation refers to track or charged-particle multiplicity, or to a summed transverse momentum value. In a given sample of events, each partial regression coefficient  $b_{ij.kl}$  denotes how a chosen statistical quantity  $x_i$  varies with one of the others,  $x_j$ , on average, for fixed values of the other two. A system of normal equations can be solved to determine these coefficients from the event sample.

The indices 1 and 2 are here used respectively to denote forward and backward particle-level quantities, and the indices 3 and 4 to denote forward and backward observed quantities, as measured in the detector. Each quantity may be correlated with any of the others. Particularly strong partial correlations exist between  $x_1$  and  $x_3$ , of course, and likewise between  $x_2$  and  $x_4$ . The regression coefficients are determined by the use of MC event samples in which both the particle-level and the simulated detector-level quantities are known.

In order to link a FB correlation at the particle level  $\rho_{\text{part}} \equiv \rho_{12}$  to the corresponding observed FB correlation  $\rho_{\text{obs}} \equiv \rho_{34}$ , we start from the second and third lines of (6.1). These equations are multiplied by  $x_3$  and by  $x_2$ , respectively, and after using the definitions of equations (2.1) and (2.2) we obtain:

$$\rho_{\text{part}} = \alpha + \beta \cdot \rho_{\text{obs}},
 \tag{6.2}$$

where

$$\alpha = \frac{b_{23.14}\sigma_3^2 - b_{32.14}\sigma_2^2 + b_{21.34}\sigma_1\sigma_3\rho_{13} - b_{34.12}\sigma_2\sigma_4\rho_{24}}{b_{31.24}\sigma_1\sigma_2}, \quad \beta = \frac{b_{24.13}\sigma_3\sigma_4}{b_{31.24}\sigma_1\sigma_2}.
 \tag{6.3}$$

The particle-level correlations are therefore linearly related to the observed correlations and may be obtained from them by means of the two coefficients  $\alpha$  and  $\beta$ .

The relationships between  $x_1$  and  $x_3$ , and between  $x_2$  and  $x_4$ , are due largely to the effect of track reconstruction efficiency; these relationships are contained in the coefficient  $\beta$ . The other terms contain detector effects such as trigger and vertex finding efficiencies, and also effects of physical FB correlations, as modelled in the MC simulations. These affect the coefficient  $\alpha$ , where they partially cancel. The values of  $\alpha$  and  $\beta$  were calculated using a range of MC tunes together with the detector simulation, and the variations on the values obtained then give rise to a systematic uncertainty on the results. This uncertainty has been reduced to a low level by this method and is contained mainly within the coefficient  $\alpha$ , whose value was found to lie in the range  $0.07 \pm 0.03$ . The value of  $\beta$  lies in the range 0.96–0.97 (0.97–0.98) for the multiplicity (summed  $p_T$ ) measurements. The full set of systematic uncertainties on the results is discussed in section 7.

To validate the above approach, the linearity of the relationship 6.2 was checked for both MC and data. It was verified from the MC samples that the analysis procedure regenerated the correlations that were calculated from the generated particles directly. The data sample did not include events produced with fewer than two reconstructed tracks associated with the primary vertex; a correction of up to 0.9% was applied for the small bias associated with this.

## 6.2 Azimuthal distributions

The azimuthal analysis employed the same track and event selection as the FB correlations analysis (section 4) but imposed a higher minimum transverse momentum requirement of  $p_{T\min}=500$  MeV on accepted tracks. This minimises systematic uncertainties by selecting a kinematic region with track reconstruction efficiencies that are flat in  $p_T$ , a procedure made necessary by the subtraction methods used in this section of the analysis. Histograms of  $N_{\text{sub}}^T$  and  $N_{\text{SO}}^T$  were made as a function of  $\Delta\phi$ , and required corrections for detector effects in order to obtain the particle-level distributions.

One set of effects comes from a loss of non-leading tracks due to detector inefficiency, and contamination from non-primary tracks. To correct for these effects, each non-leading track was weighted by a factor  $(1-b)/\epsilon$ , where  $b$  represents the fraction of background due to non-primary tracks and  $\epsilon$  is the mean track reconstruction efficiency for the selected tracks, evaluated as a function of  $p_T$  and  $\eta$ . These terms were evaluated from MC simulation. The normalisation procedure gives a large cancellation of the effects of non-leading track reconstruction inefficiency, while preserving the shape of the variations in  $\epsilon$ ; these are more important than the average value.

The loss of a leading track, due to track reconstruction inefficiency, may lead to the loss of the event or else to the wrong identification of a non-leading track as the leading track. A correction was applied to the observed distributions of figure 1 to remove the shape distortions caused by these effects. It was done in three steps. First, in a given  $p_T$  and  $\eta$  range, the fraction  $\epsilon_l$  of events that have the correct leading track was estimated by integrating the track reconstruction efficiency (determined from MC-simulated events) over the observed  $p_T$  and  $\eta$  spectrum of the events in the data. Second, starting from the observed data distribution in  $\Delta\phi$ , a series of new distributions was constructed by removing an increasingly large random fraction  $f$  of the observed leading tracks from the

data sample. The resulting shape distortion was found to be linear in  $f$  in each  $\Delta\phi$  bin. This made it possible to extrapolate the original data distribution to what it would have been if  $\epsilon_l$  were 100%. The correction to the final results is up to 2%.

The distributions of  $N_{\text{sub}}^T$  and  $N_{\text{SO}}^T$  were made for three ranges of  $\eta$ , namely  $|\eta| < 1.0$ , 2.0 and 2.5, in order to provide a broad-based comparison with the MC models.

## 7 Systematic uncertainties

### 7.1 Forward-backward correlations

The largest systematic uncertainty on the results arises from the uncertainty on the track reconstruction efficiency. A smaller effect arises from the variation of the correction parameters  $\alpha$  and  $\beta$  with the MC model used to obtain them (cf. section 6.1). The effects from the uncertainties on trigger and vertex finding efficiencies are small.

In this section, uncertainties are given as percentages of the central values and symmetric positive and negative uncertainties are taken. The model-dependence of the analysis was evaluated for the 7 TeV data, and the conclusions were also applied to the 900 GeV data. The chosen initial set of five MC tune samples were used to correct the data, and the spread in the results relative to the standard MC09, of the order of 1%, was taken as the systematic uncertainty due to the model-dependence of the correction procedure. In addition, the diffractive components of the MC09 sample were varied by  $\pm 10\%$ ; this changed the measured correlations by typically  $\pm 0.8\%$  of their value.

The systematic effects on the FB correlations due to uncertainties on the trigger, vertex finding and track reconstruction efficiencies were studied by varying these efficiencies in the MC, and then re-evaluating the FB correlations using the standard correction parameters. Uncertainties on efficiencies of 0.9% (1.6%) and 0.4% (0.6%) were used at 7 TeV (900 GeV) for the trigger and vertex finding, respectively, and a track reconstruction uncertainty varying between 2% and 15% [17] was used. The systematic variations on the standard selection were achieved by removing a fraction of the events or tracks in a sample, randomly selected, according to the uncertainties on the respective efficiencies. This simulated a decrease in the efficiency which was taken as the corresponding systematic uncertainty.

In the case of the trigger and vertex efficiencies, the deviations on the FB correlations are  $\leq 0.2\%$  and  $\sim 0.2\%$ , respectively. When the number of tracks in the events was varied as described above, the resulting FB correlation varied by 1.2% (1.4%) in the central pair of  $\eta$ -intervals, increasing to 3.3% (3.0%) in the outermost pair for multiplicity (momentum).

The total systematic uncertainty was obtained by summing the individual contributions in quadrature, and was combined in quadrature with the statistical uncertainty to obtain the total uncertainty. The total systematic uncertainty is dominated by the errors on the model-dependence and on the track reconstruction efficiency, and both of these are highly correlated between the two energies. In taking the ratios of the correlations at different energies, the systematic uncertainties are therefore small relative to the statistical and can be neglected.

## 7.2 Azimuthal distributions

Performing the subtractions and the normalisations reduces the systematic uncertainties on the  $\Delta\phi$  distributions considerably. Of primary importance is the uncertainty on the shape of the distribution, quantified here by estimating the relative uncertainty on the first bin in  $\Delta\phi$ . The most important effect comes from the event selection efficiency (1% in the 900 GeV data and 3% in the 7 TeV data), and mainly affects events with a small number of tracks. Other important effects come from the correction procedure, which contributes 2%, and the effects of the track momentum resolution in the vicinity of the boundaries of the kinematic acceptance. The latter was estimated by varying the values of the cuts on  $p_T$  and  $\eta$  by their resolution and repeating the analysis. This gave variations of 1–2% and 0.2–0.5% from  $p_T$  and  $\eta$  respectively.

A variety of additional effects gave uncertainties at the 0.1–0.3% level. These comprised errors in the choice of leading charged particle (estimated from MC), uncertainty on the fraction of leading particles that were not reconstructed (also estimated from MC), and uncertainties on the track reconstruction efficiencies of the non-leading tracks. The uncertainty on the track reconstruction efficiency, measured to be on average 15% [17], was propagated to the non-leading track correction. It corresponds to a relative systematic uncertainty on the distribution of 0.1–0.2%. A final contribution of similar magnitude came from uncertainties on the number of secondary particles, depending on the choice of impact parameter cut. The track reconstruction efficiency difference between the forward and backward halves of the detector was found to have a negligible effect.

The total systematic uncertainty was obtained by summing the individual contributions in quadrature, and was combined in quadrature with the statistical uncertainty to obtain the total uncertainty. The statistical uncertainties dominate the measurements at 900 GeV, but in the 7 TeV data the statistical and systematic uncertainties are comparable.

## 8 Results

### 8.1 Forward-backward correlations

The observed FB correlations were corrected to the particle level, using the procedure described in section 6, and all results are quoted at this level. The correction typically increased the correlation value above the observed, uncorrected value by 0.05–0.08, in both multiplicity and momentum.

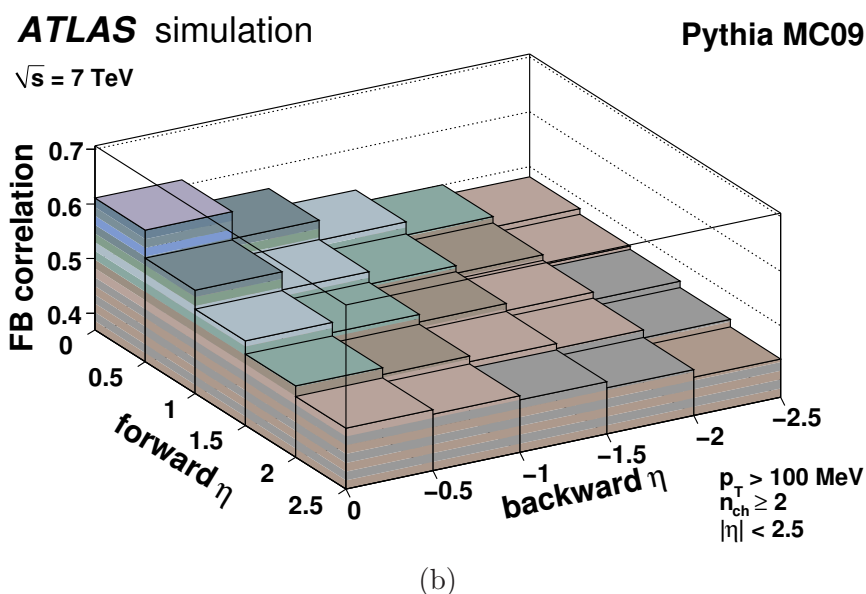
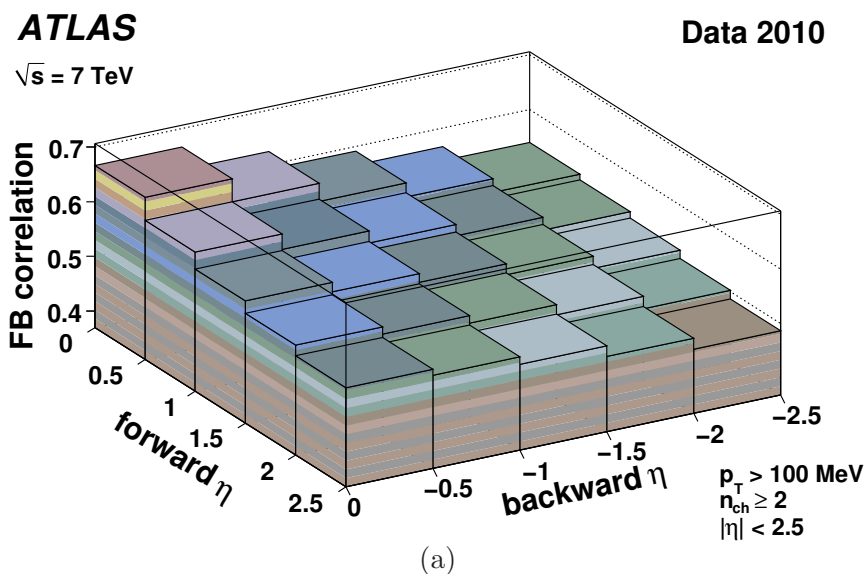
To obtain an overall picture, and as a check on the experimental approach, the FB multiplicity correlations using 7 TeV collisions were calculated over a matrix of forward-backward pairs of  $\eta$  intervals of width  $\delta\eta = 0.5$  covering the range 0 to 2.5 in each direction, for a  $p_{T\min}$  value of 100 MeV. The results are given in table 1 and shown pictorially in figure 2. The main diagonal depicts the symmetric FB correlations. Both for the data and for the chosen MC, the correlations are seen to vary strongly with the separation between the chosen intervals, and only weakly with the mean  $\eta$  value for a given separation, confirming that the symmetric forward-backward configuration is a characteristic representation of the correlation at a given separation in  $\eta$ .

Forward $\eta$ interval Backward $\eta$ interval	0.0–0.5	0.5–1.0	1.0–1.5	1.5–2.0	2.0–2.5
0.0–0.5	0.666 $\pm 0.011$	0.624 $\pm 0.011$	0.592 $\pm 0.011$	0.566 $\pm 0.012$	0.540 $\pm 0.013$
0.5–1.0	0.624 $\pm 0.011$	0.596 $\pm 0.011$	0.574 $\pm 0.012$	0.553 $\pm 0.013$	0.530 $\pm 0.014$
1.0–1.5	0.594 $\pm 0.011$	0.576 $\pm 0.012$	0.560 $\pm 0.013$	0.540 $\pm 0.014$	0.518 $\pm 0.014$
1.5–2.0	0.571 $\pm 0.012$	0.557 $\pm 0.013$	0.544 $\pm 0.014$	0.526 $\pm 0.014$	0.503 $\pm 0.016$
2.0–2.5	0.551 $\pm 0.013$	0.540 $\pm 0.014$	0.527 $\pm 0.014$	0.507 $\pm 0.016$	0.487 $\pm 0.018$

**Table 1.** Multiplicity correlations for events at  $\sqrt{s} = 7$  TeV for events with a minimum of two charged particles in the kinematic interval  $p_T > 100$  MeV and  $|\eta| < 2.5$  for different combinations of forward and backward pseudorapidity interval. The quoted uncertainties are the systematic uncertainties; the statistical uncertainties are approximately  $\pm 0.001$  or less and are small by comparison.

In the remainder of this section, the correlations between symmetrically opposite FB pseudorapidity intervals of width  $\delta\eta = 0.5$  are examined. A global correlation has also been calculated, defined as the FB correlation taken over the entire pseudorapidity range ( $\delta\eta = 2.5$ ). The results are summarised in table 2. Figure 3(a) shows the variation of the FB multiplicity correlation with  $\eta$  found in the 7 TeV data, compared with various MC tunes. Figure 3(b) shows the same quantities measured at 900 GeV. The general shape of the data distributions is reasonably well reproduced by the MCs, though there are discrepancies in detail. In the case of MC09, the trend is to underestimate the correlation by  $\sim 10\%$ . The correlation is generally similar for the DW and MC09 tunes at 7 TeV, but shows greater variation at 900 GeV where it rises relative to the data at high  $\eta$  values, implying an inaccuracy in the modelled ratio of long-distance to short-distance correlations. The PYTHIA 8 tune performs well at low  $\eta$  but somewhat poorly at higher values. The correlations given by AMBT2B are the most consistent with the data; AMBT2B was tuned to these event samples using different variables. At both energies HERWIG++ gives correlations consistently higher than the data.

Figure 3(c) shows the ratio of the FB multiplicity correlation distributions at 900 GeV and 7 TeV, compared with the results from the MC tunes. The multiplicity correlations are substantially lower at 900 GeV than at 7 TeV and the relative difference is greater for the larger pseudorapidity bins. This is a new experimental result which can be interpreted in general terms as indicating that, while the short-range correlation is similar at the two energies, the long-range correlation becomes considerably higher as the energy increases. This trend is reproduced in most of the MC predictions, although not so well by DW. The HERWIG++ ratio is not included since different tunes are used at the two energies.



**Figure 2.** Multiplicity correlation at 7 TeV in a matrix of forward/backward  $\eta$ -intervals for  $|\eta| < 2.5$ , for events with at least two charged particles with  $p_T > 100 \text{ MeV}$  and  $|\eta| < 2.5$ . (a) data; (b) PYTHIA 6 MC09 tune. The statistical uncertainties are of the order of  $\pm 0.001$ .

The minimum value,  $p_{T\text{min}}$ , of the transverse momentum of the selected charged particles was varied for the 7 TeV data. The resulting multiplicity correlations are shown in table 3 and illustrated in figure 4. As expected, the correlations fall rapidly as  $p_{T\text{min}}$  increases above a few hundred MeV, a feature also seen in the MC models (not shown). In the context of the Lund string model [9], at low transverse momentum values, this may be seen as a general tendency for a partonic string to fragment in a uniform way all along its length. There is also a possible influence of MPI effects. At higher transverse momentum, however, particles are more likely to be associated with jets, and there is no strong corre-

$ \eta $ interval	0.0 – 0.5	0.5 – 1.0	1.0 – 1.5	1.5 – 2.0	2.0 – 2.5	0.0 – 2.5
900 GeV	0.498 $\pm 0.001 \pm 0.011$	0.407 $\pm 0.001 \pm 0.011$	0.359 $\pm 0.002 \pm 0.013$	0.310 $\pm 0.002 \pm 0.015$	0.259 $\pm 0.002 \pm 0.015$	0.652 $\pm 0.001 \pm 0.007$
7 TeV	0.666 $\pm 0.011$	0.596 $\pm 0.011$	0.560 $\pm 0.013$	0.526 $\pm 0.014$	0.487 $\pm 0.018$	0.801 $\pm 0.008$

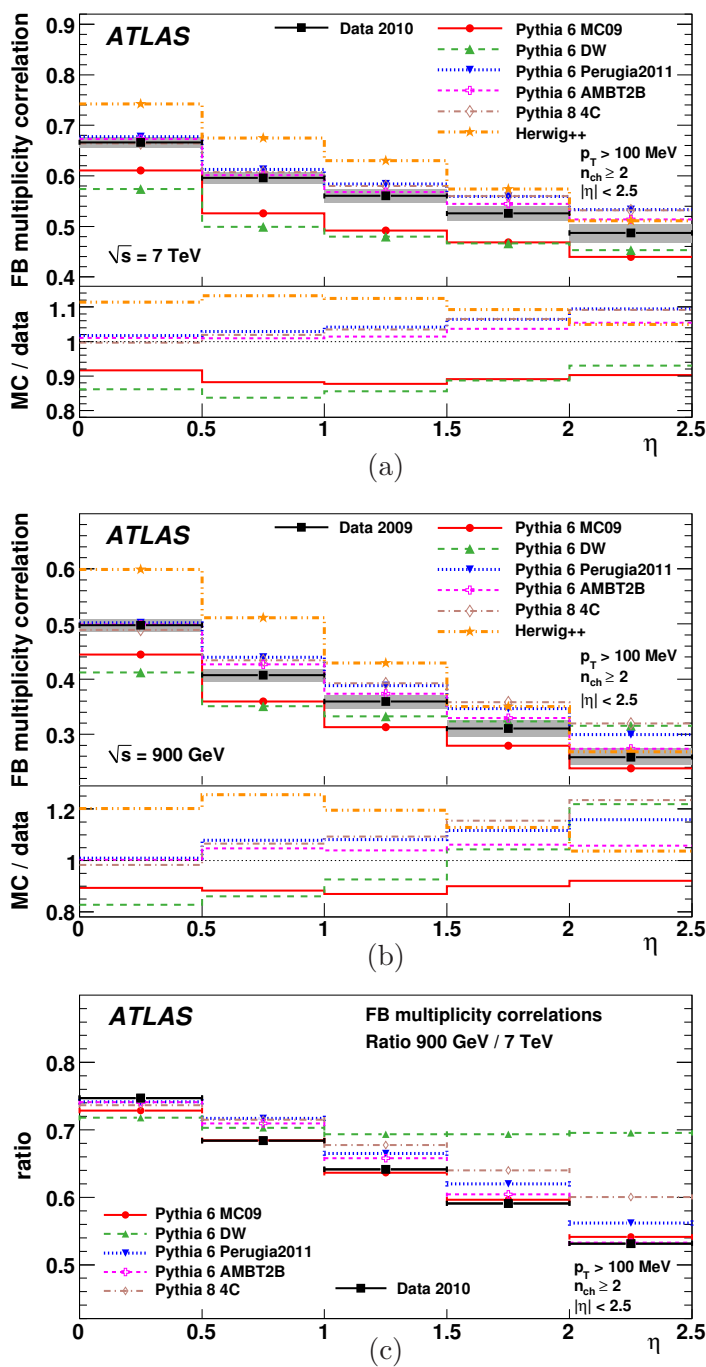
**Table 2.** Forward-backward multiplicity correlation for charged particles in symmetrically opposite  $\eta$ -intervals for events with a minimum of two charged particles in the kinematic interval  $p_T > 100$  MeV and  $|\eta| < 2.5$  at  $\sqrt{s} = 900$  GeV and 7 TeV. The final column shows the global result. The quoted uncertainties are first statistical, second systematic; where the statistical uncertainty is less than approximately  $\pm 0.001$ , only the systematic uncertainty is shown.

$ \eta $ interval $p_{T\min}$ [GeV]	0.0 – 0.5	0.5 – 1.0	1.0 – 1.5	1.5 – 2.0	2.0 – 2.5	0.0 – 2.5
0.1	0.666 $\pm 0.011$	0.596 $\pm 0.011$	0.560 $\pm 0.013$	0.526 $\pm 0.014$	0.487 $\pm 0.018$	0.801 $\pm 0.008$
0.2	0.639 $\pm 0.009$	0.565 $\pm 0.010$	0.534 $\pm 0.011$	0.498 $\pm 0.012$	0.458 $\pm 0.017$	0.790 $\pm 0.008$
0.3	0.602 $\pm 0.009$	0.529 $\pm 0.009$	0.500 $\pm 0.011$	0.462 $\pm 0.012$	0.426 $\pm 0.016$	0.775 $\pm 0.008$
0.5	0.518 $\pm 0.009$	0.445 $\pm 0.009$	0.416 $\pm 0.010$	0.377 $\pm 0.014$	0.341 $\pm 0.017$	0.726 $\pm 0.010$
1.0	0.291 $\pm 0.010$	0.216 $\pm 0.011$	0.190 $\pm 0.013$	0.158 $\pm 0.018$	0.132 $\pm 0.017$	0.506 $\pm 0.023$
1.5	0.128 $\pm 0.001 \pm 0.006$	0.063 $\pm 0.001 \pm 0.008$	0.047 $\pm 0.001 \pm 0.016$	0.026 $\pm 0.001 \pm 0.012$	0.010 $\pm 0.001 \pm 0.014$	0.234 $\pm 0.033$
2.0	0.036 $\pm 0.001 \pm 0.006$	-0.020 $\pm 0.001 \pm 0.008$	-0.031 $\pm 0.001 \pm 0.007$	-0.041 $\pm 0.001 \pm 0.005$	-0.046 $\pm 0.001 \pm 0.007$	-0.008 $\pm 0.001 \pm 0.019$

**Table 3.** Forward-backward charged-particle multiplicity correlations at  $\sqrt{s} = 7$  TeV, in symmetrically opposite  $\eta$ -intervals, for events with at least two charged particles with  $|\eta| < 2.5$  and  $p_T > p_{T\min}$ , for varying values of  $p_{T\min}$ . The final column shows the global result. The quoted uncertainties are first statistical, second systematic; where the statistical uncertainty is less than approximately  $\pm 0.001$ , only the systematic uncertainty is shown.

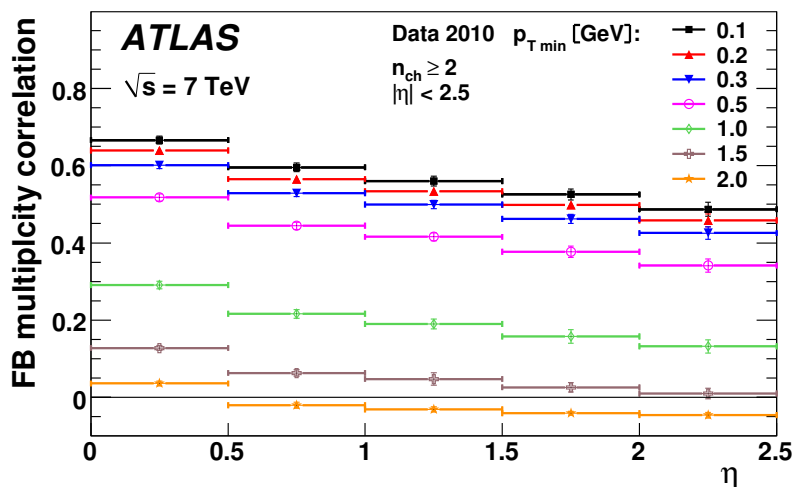
lation between a given jet and another jet at any particular value of  $\eta$ . The tendency to have negative correlation values at higher  $p_T$  is a manifestation of the tendency of high- $p_T$  particles to be collimated into jets.

The FB correlations between the summed  $p_T$  of charged particles are shown in figure 5 and listed in table 4. In general, the features are similar to those of the multiplicity correlations. As in the case of the multiplicities, the momentum FB correlations are higher at 7 TeV than at 900 GeV.

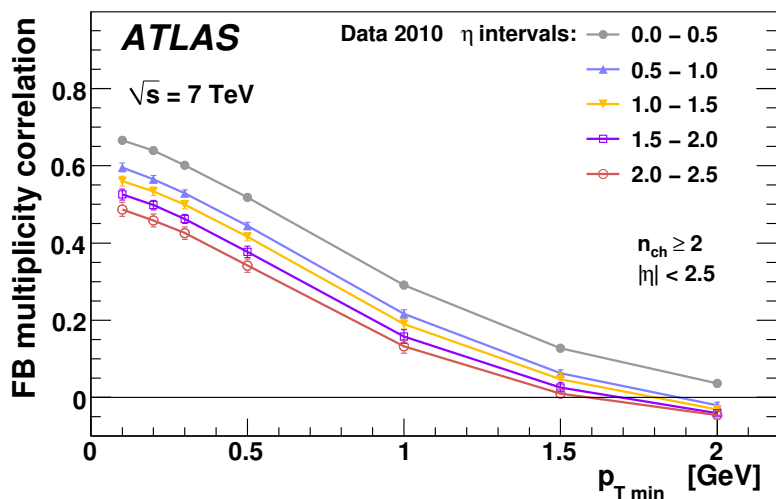


**Figure 3.** Forward-backward multiplicity correlation in symmetrically opposite  $\eta$  intervals for events with at least two charged particles with  $p_T \geq 100$  MeV and  $|\eta| < 2.5$ . (a) Data at 7 TeV, compared with a selection of MC simulations. The systematic uncertainties are indicated by a grey band; the statistical uncertainties are too small to be visible on the figure. (b) The same at 900 GeV. (c) Ratio of the 900 GeV results to the 7 TeV results.





(a)



(b)

**Figure 4.** Forward-backward multiplicity correlations: (a) as a function of  $\eta$  for different values of  $p_{T\min}$ ; (b) as a function of  $p_{T\min}$  for different values of  $\eta$ , with lines drawn to guide the eye. The plotted systematic uncertainties are smaller than the symbol in many cases; the statistical uncertainties are too small to be visible on the figure.

## 8.2 Azimuthal distributions

Distributions of the subtracted and the “same hemisphere minus opposite” charged-particle distributions in azimuthal angle,  $N_{\text{sub}}^T$  and  $N_{\text{SO}}^T$ , are shown in figures 6 and 7, respectively, for  $\sqrt{s} = 900 \text{ GeV}$  and  $7 \text{ TeV}$ , and compared with a selection of the different MC tunes. The leading particle in each event is omitted from the plots. One notes in figure 6 the increasing prominence of the “opposite” jet peak in the distributions as the  $|\eta|$  interval increases, as more pairs of jets, separated in  $\eta$ , become included within the acceptance. The shapes show a strong energy dependence.

$ \eta $ interval	0.0 – 0.5	0.5 – 1.0	1.0 – 1.5	1.5 – 2.0	2.0 – 2.5	0.0 – 2.5
900 GeV	0.445 $\pm 0.001 \pm 0.008$	0.362 $\pm 0.002 \pm 0.009$	0.328 $\pm 0.002 \pm 0.010$	0.284 $\pm 0.002 \pm 0.013$	0.236 $\pm 0.003 \pm 0.009$	0.637 $\pm 0.001 \pm 0.009$
7 TeV	0.590 $\pm 0.010$	0.519 $\pm 0.009$	0.492 $\pm 0.010$	0.455 $\pm 0.011$	0.412 $\pm 0.016$	0.770 $^{+0.007}_{-0.008}$

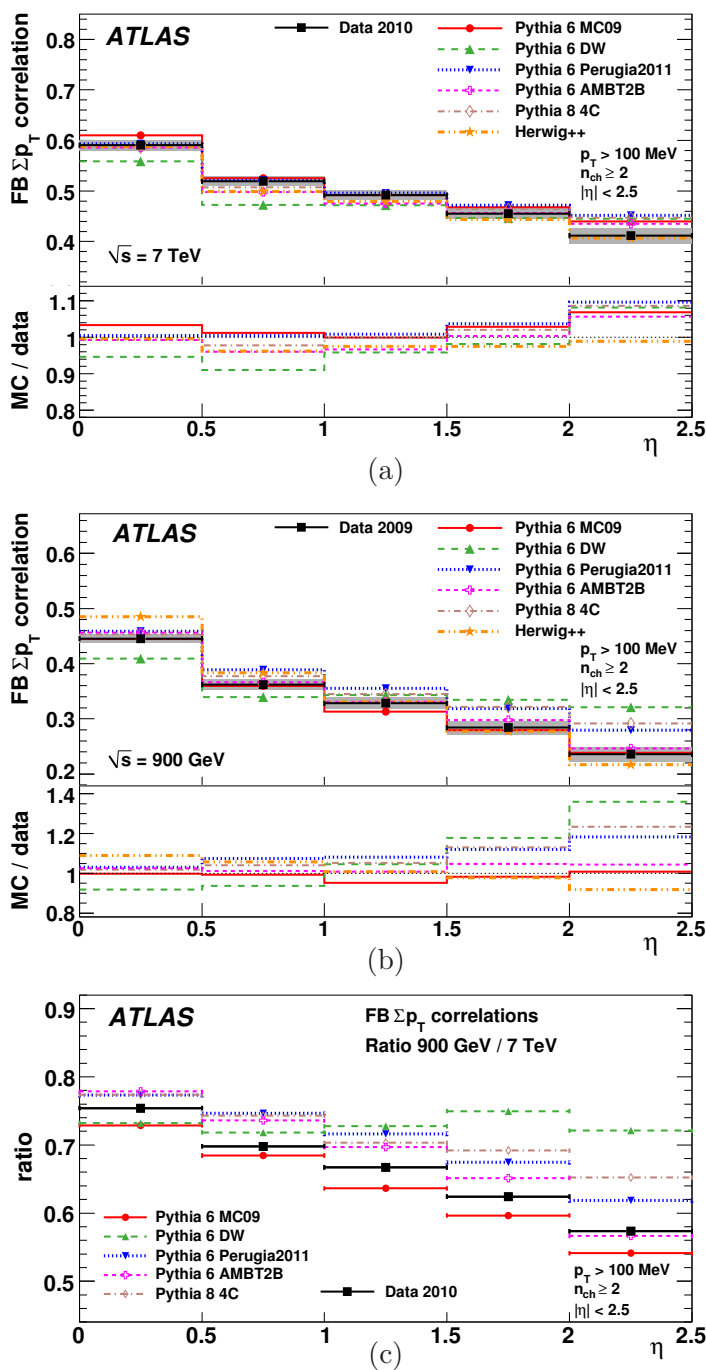
**Table 4.** Forward-backward summed  $p_T$  correlation for charged particles in symmetrically opposite  $\eta$  intervals, for events with at least two charged particles in the kinematic interval  $p_T > 100$  MeV and  $|\eta| < 2.5$  at  $\sqrt{s} = 900$  GeV and 7 TeV, for corrected data. The quoted uncertainties are first statistical, second systematic; where the statistical uncertainty less than approximately  $\pm 0.001$ , only the systematic uncertainty is shown.

In general, the MC tunes and models used for comparison with data provide a reasonable description in the restricted rapidity range  $|\eta| < 1$ , but the agreement becomes increasingly poor as the  $|\eta|$  range is increased to 2 and 2.5. For the subtracted distributions the Perugia2011 tune is the most successful, giving a good representation of the distribution at 900 GeV for all three  $|\eta|$  ranges. At 7 TeV it is poorer for  $|\eta| < 1$ , but similar to the other tunes and models presented, and it improves as the coverage in  $|\eta|$  increases. The other three models are very similar at 900 GeV, where they all give an equally poor description of the data in the broader  $|\eta|$  ranges. At 7 TeV, DW describes the data well for  $|\eta| < 1$ , but is a poorer match in the higher  $|\eta|$  ranges. HERWIG++ and PYTHIA 8 are similar at all  $|\eta|$ , giving a somewhat poor description that resembles Perugia2011 at  $|\eta| < 1$  and DW at higher  $|\eta|$ . In all cases, the discrepancies of the tunes and models tend towards over-enhancing the same-side peak relative to the opposite-side peak.

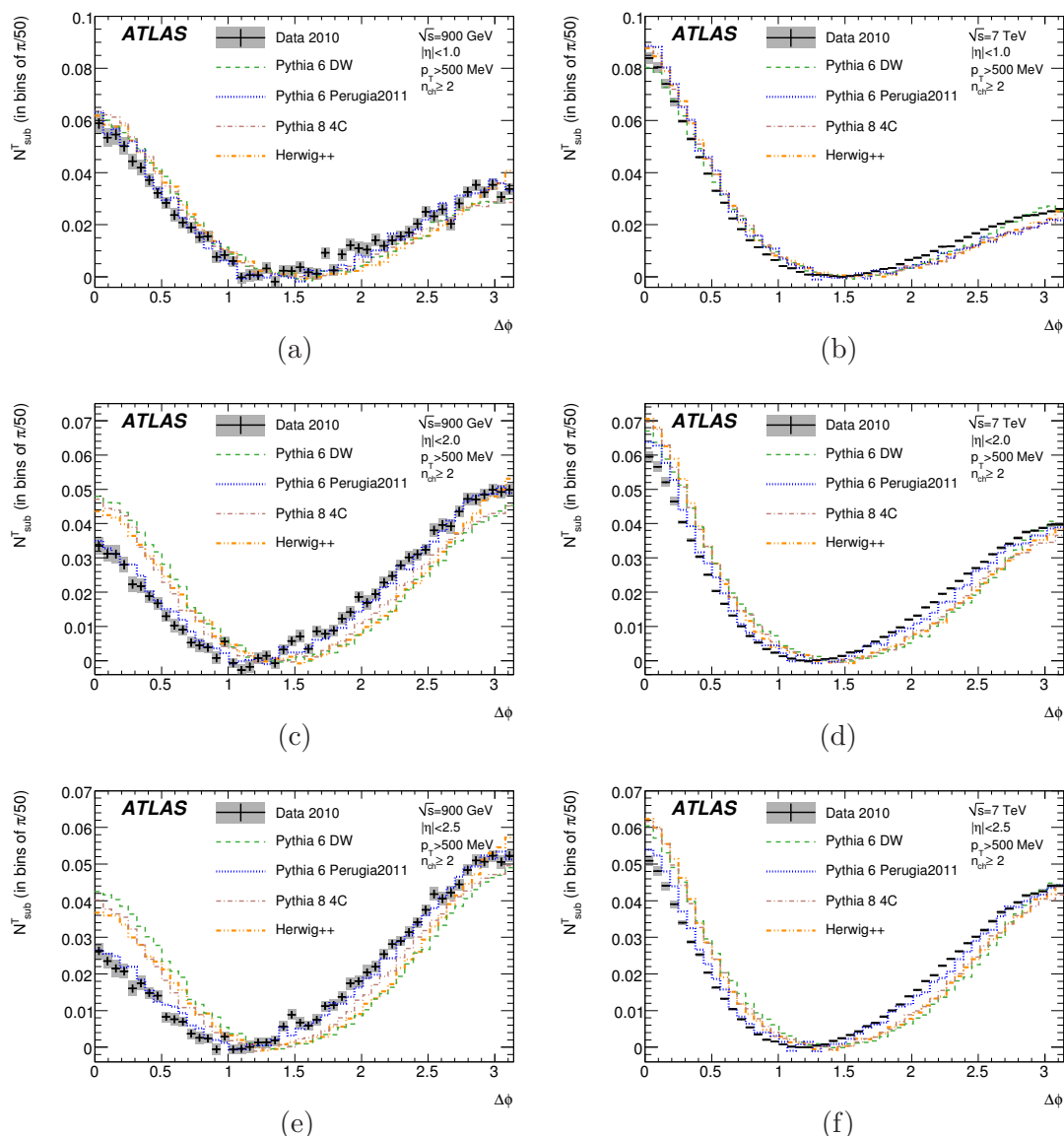
The same-minus-opposite variable is described well for  $|\eta| < 1$  at both energies by PYTHIA 8, Perugia2011 and HERWIG++, with regard to both the leading-particle “jet peak” and the opposing pedestal. However DW exaggerates the sharpness of the peak. As the accepted range in  $|\eta|$  increases, so that the averaged effects of more widely-separated jet pairs become included in the distributions, the disagreements become more pronounced. At both energies DW’s exaggeration of the peak increases dramatically with  $|\eta|$ , while PYTHIA 8 gives a progressively worse description of the peak, also overestimating its size. The HERWIG++ and Perugia2011 models are very similar in this variable, and describe the data well at 900 GeV but less so at 7 TeV, where they also resemble PYTHIA 8. Again, the discrepancies tend towards an over-enhancement of the same-side “jet peak”, which becomes a particularly strong effect in the case of DW.

## 9 Conclusions

The forward-backward correlation in charged-particle multiplicity and in summed charged-particle  $p_T$  has been measured at ATLAS in minimum-bias events at  $\sqrt{s} = 900$  GeV and 7 TeV. This is the first measurement of summed- $p_T$  correlations. Both types of correlation fall in strength with increasing separation in pseudorapidity. The multiplicity correlation is found to differ by up to about 15% from some of the standard Monte Carlo tunes at



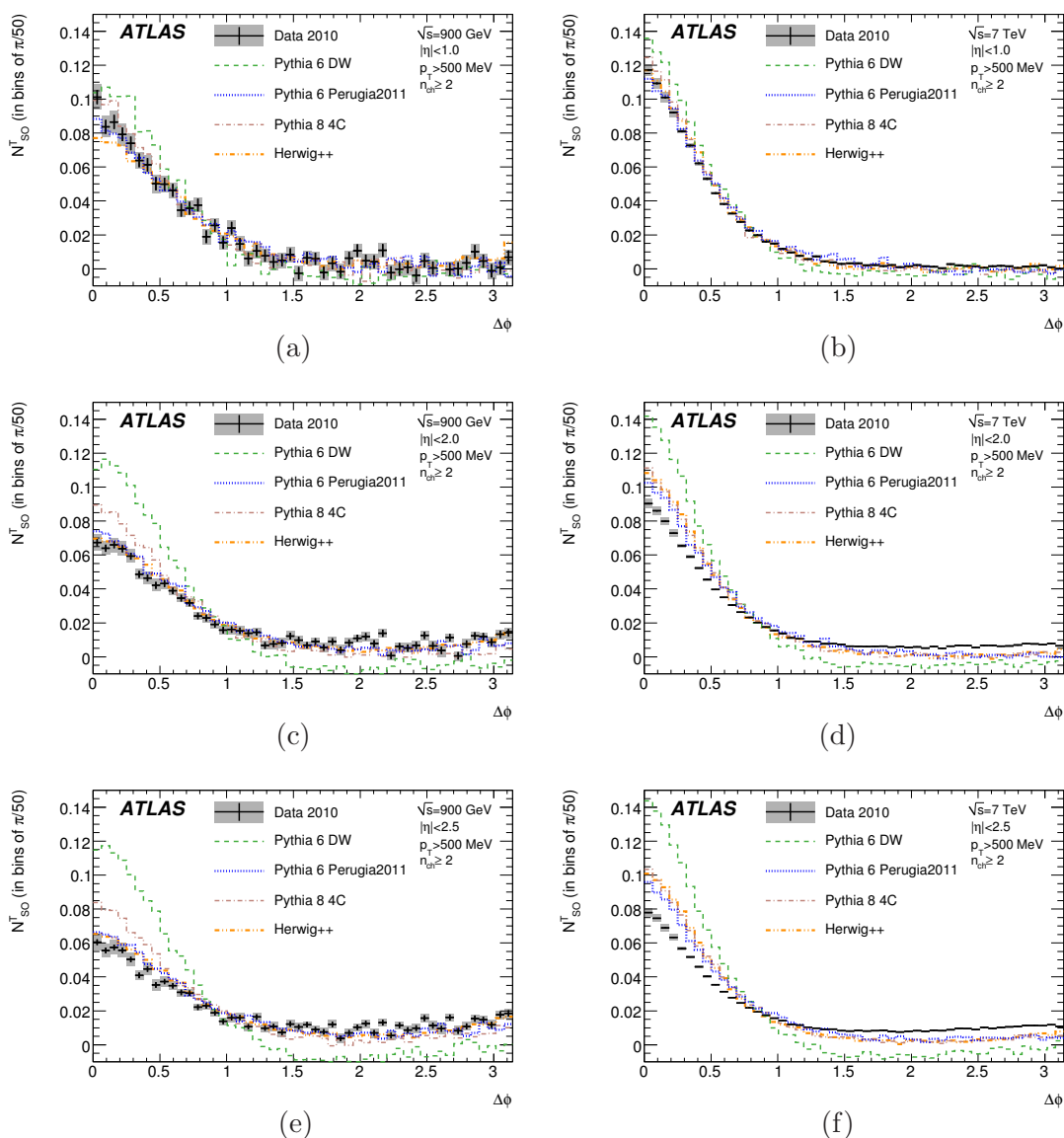
**Figure 5.** Forward-backward  $\Sigma p_T$  correlation in symmetrically opposite  $\eta$  intervals for events with at least two charged particles with  $p_T \geq 100 \text{ MeV}$  and  $|\eta| < 2.5$ . (a) Data at 7 TeV, compared with a selection of MC simulations. The systematic uncertainties are indicated by a grey band; the statistical uncertainties are too small to be visible on the figure. (b) The same at 900 GeV. (c) Ratio of the 900 GeV results to the 7 TeV results.



**Figure 6.** Subtracted and normalised  $\Delta\phi$  distributions ( $N_{\text{sub}}^T$ ) for data taken at  $\sqrt{s} = 900$  GeV (a, c, e) and 7 TeV (b, d, f). Three different pseudorapidity regions are shown:  $|\eta| < 1.0$ ,  $|\eta| < 2.0$ ,  $|\eta| < 2.5$ . The data (black symbols) are compared with the predictions of different MC tunes (see text for more details). The vertical bars on the data points denote the statistical uncertainty, while the shaded areas denote the total uncertainty. The leading particle in each event is omitted from the plot.

7 TeV, and by somewhat more at 900 GeV; the shapes of the distributions are generally similar but some of the tunes show systematically different trends from the data. With the summed- $p_T$  correlation, similar conclusions are found. It should be noted that present Monte Carlo generators have not been tuned using these correlations.

Both types of correlation are found to be significantly stronger at 7 TeV than at 900 GeV, the relative increase being greater at larger intervals in pseudorapidity. This



**Figure 7.** Normalised “same hemisphere minus opposite” distributions ( $N_{SO}^T$ ) for data taken at  $\sqrt{s} = 900$  GeV (a, c, e) and 7 TeV (b, d, f). Three different pseudorapidity regions are shown:  $|\eta| < 1.0$ ,  $|\eta| < 2.0$ ,  $|\eta| < 2.5$ . The data (black symbols) are compared with the predictions of different MC tunes (see text for more details). The vertical bars on the data points denote the statistical uncertainty, while the shaded areas denote the total uncertainty. The leading particle in each event is omitted from the plot.

trend is reproduced in most of the MC models. As expected, a strong tendency is seen for the multiplicity correlations to fall with increasing minimum transverse momentum of the charged particles. This illustrates well the transition between the soft, non-perturbative regime of parton string or cluster fragmentation and the jet-dominated regime of perturbative quantum chromodynamics.

Azimuthal charged-particle distributions relative to the leading charged particle have

been measured by means of two variables, using the full pseudorapidity reach of the ATLAS tracking detectors. These variables minimise inter-jet effects in different ways, and give different sensitivities to the jet-like properties of the event distributions. In both cases, the MC models give fair to good descriptions of the data, provided that the selected pseudorapidity range is within one unit of zero. As the range in pseudorapidity is increased, however, disagreements with most of the models become more pronounced, in describing both the jet-influenced region near the leading charged particle and also the opposite-side region. This indicates the complexity of the effects that the models are seeking to describe when broad features of the events are being considered.

The event features studied here provide new means of discriminating between the different possible tunes of the Monte Carlo models. They will need to be taken into account, if the most accurate description of particle processes at the highest LHC energies is to be attained.

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## A Summary of Monte Carlo tunes

Table 5 summarises the main features of the PYTHIA- and HERWIG-based Monte Carlo tunes employed in the analysis and referred to in the text.

Generator	Tune	Shower ordering	PDF	Additional information
PYTHIA 6.4	DW	virtuality	CTEQ 5L	Non-interleaved MPI model
	MC09	$p_T$	MRSTLO*	MPI interleaved with ISR
	Perugia 2011	$p_T$	CTEQ 5L	MPI interleaved with ISR
	AMBT2B	$p_T$	CTEQ 6L1	MPI interleaved with ISR
PYTHIA 8.150	4C	$p_T$	CTEQ 6L1	MPI interleaved with ISR and FSR; updated model for diffraction
HERWIG++ 2.5	MU900-2	angle	MRST 2007LO*	MB tune for 900 GeV only; no model for diffraction
	UE7-2	angle	MRST 2007LO*	UE tune for 7 TeV only; no model for diffraction

**Table 5.** Main features of the different PYTHIA and HERWIG tunes used in the comparisons with the data, as used with CTEQ [37, 38] and MRST [39] parton distribution functions.

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